

### 3 Dynamical mass

(idea by A. Milekhin)

A free massive scalar field  $\phi$  in 1 + 1 dimensions is described by

$$\mathcal{L} = (\partial_t \phi)^2 - (\partial_x \phi)^2 - m^2 \phi^2$$

In many physical problems it is useful to promote coupling constants to dynamical fields. We can do the same with the mass  $m$  considering the following partition function:

$$Z = \int \mathcal{D}\phi \mathcal{D}m \exp \left( i \int dt dx (\partial_t \phi)^2 - (\partial_x \phi)^2 - m(x)^2 \phi^2 \right)$$

Study the stability of the following vacuum:  $\phi = 0$ ,  $m = m_0 = \text{const} \neq 0$ . Consider two cases: infinite volume and a finite interval  $[0, L]$  in the  $x$  direction with Dirichlet boundary conditions for  $\phi$  and free boundary conditions for  $m$ .

#### Solution:

To check if the configuration  $\phi = 0, m = m_0 = \text{const}$  is indeed vacuum, first we need to assume that at infinity  $\phi \rightarrow \phi_{cl}$  and  $m$  is constant. Then we will investigate the form of the effective potential for  $\phi_{cl}$  and  $m$ . Integrating out  $\phi$  by computing  $\det(\partial^2 + m^2)$  produces the celebrated Coleman–Weinberg effective potential:

$$V_{eff}(m, \phi_{cl}) = m^2 \log \frac{m^2}{M^2} - m^2 + m^2 \phi_{cl}^2 \tag{1}$$

where  $M$  is the renormalization scale. The potential has minimum at  $m = 0$  and  $\phi_{cl} = -\infty$ . However, at  $m = 0$  the theory has symmetry  $\phi \rightarrow \phi + \text{const}$ , therefore  $\phi_{cl}$  can be renormalized to any value. For example to  $\phi_{cl} = 0$ .

On the infinite line because of the translation invariance it is obvious that configuration with constant  $m$  and  $\phi$  is stable under non-uniform perturbations.

However, on a finite interval with Dirichlet boundary conditions the translation invariance is broken. It turned out that it is quite hard to find the exact  $m(x)$  profile, however it is easy to show that  $m(x) = \text{const}$  is not even a saddle-point of the effective potential. We refer to the original paper [1](see also [2]) for the derivation.

#### References

- [1] S. Bolognesi, K. Konishi and K. Ohashi, JHEP **1610**, 073 (2016) doi:10.1007/JHEP10(2016)073 [arXiv:1604.05630 [hep-th]].
- [2] A. Milekhin, Phys. Rev. D **95**, no. 8, 085021 (2017) doi:10.1103/PhysRevD.95.085021 [arXiv:1612.02075 [hep-th]].